

No-Signalling Is Equivalent To Free Choice Of Measurements

Samson Abramsky, Adam Brandenburger and Andrei Savochnik

Department of Computer Science
The University of Oxford

Introduction

Introduction

We shall show an equivalence between two conditions on probabilistic physical theories which are usually taken to be quite different in nature:

Introduction

We shall show an equivalence between two conditions on probabilistic physical theories which are usually taken to be quite different in nature:

- **No-Signalling**, which is usually taken to reflect constraints imposed by Relativity.

Introduction

We shall show an equivalence between two conditions on probabilistic physical theories which are usually taken to be quite different in nature:

- **No-Signalling**, which is usually taken to reflect constraints imposed by Relativity.
- **Free choice of measurements**, which seems to be a necessary assumption for any No-Go theorem.

Introduction

We shall show an equivalence between two conditions on probabilistic physical theories which are usually taken to be quite different in nature:

- **No-Signalling**, which is usually taken to reflect constraints imposed by Relativity.
- **Free choice of measurements**, which seems to be a necessary assumption for any No-Go theorem.

Our equivalence is mediated by **negative probabilities**, and is based on a striking result characterising the correlations which can be realised by local hidden-variable theories if we allow signed measures (negative probabilities) on the hidden variables.

Introduction

We shall show an equivalence between two conditions on probabilistic physical theories which are usually taken to be quite different in nature:

- **No-Signalling**, which is usually taken to reflect constraints imposed by Relativity.
- **Free choice of measurements**, which seems to be a necessary assumption for any No-Go theorem.

Our equivalence is mediated by **negative probabilities**, and is based on a striking result characterising the correlations which can be realised by local hidden-variable theories if we allow signed measures (negative probabilities) on the hidden variables.

Negative probabilities have been considered by Wigner, Dirac, Feynman *et al.* without ever acquiring a clear status.

Introduction

We shall show an equivalence between two conditions on probabilistic physical theories which are usually taken to be quite different in nature:

- **No-Signalling**, which is usually taken to reflect constraints imposed by Relativity.
- **Free choice of measurements**, which seems to be a necessary assumption for any No-Go theorem.

Our equivalence is mediated by **negative probabilities**, and is based on a striking result characterising the correlations which can be realised by local hidden-variable theories if we allow signed measures (negative probabilities) on the hidden variables.

Negative probabilities have been considered by Wigner, Dirac, Feynman *et al.* without ever acquiring a clear status.

We shall also sketch an ‘operational interpretation’ of negative probabilities.

No-Signalling

No-Signalling

The usual story:

When Alice and Bob perform measurements in spacelike-separated locations, the marginal probabilities for Alice's observations of outcomes for her measurements are independent of Bob's choice of measurement setting.

No-Signalling

The usual story:

When Alice and Bob perform measurements in spacelike-separated locations, the marginal probabilities for Alice's observations of outcomes for her measurements are independent of Bob's choice of measurement setting.

This can be taken to reflect the inability, under relativistic constraints, for information about Bob's settings to reach Alice's site in time to influence her outcomes.

No-Signalling

The usual story:

When Alice and Bob perform measurements in spacelike-separated locations, the marginal probabilities for Alice's observations of outcomes for her measurements are independent of Bob's choice of measurement setting.

This can be taken to reflect the inability, under relativistic constraints, for information about Bob's settings to reach Alice's site in time to influence her outcomes.

Reasons to be doubtful:

No-Signalling

The usual story:

When Alice and Bob perform measurements in spacelike-separated locations, the marginal probabilities for Alice's observations of outcomes for her measurements are independent of Bob's choice of measurement setting.

This can be taken to reflect the inability, under relativistic constraints, for information about Bob's settings to reach Alice's site in time to influence her outcomes.

Reasons to be doubtful:

- Probabilities refer to **ensembles**. The entire ensembles must be space-like separated from each other; and they may be unbounded in size ...

No-Signalling

The usual story:

When Alice and Bob perform measurements in spacelike-separated locations, the marginal probabilities for Alice's observations of outcomes for her measurements are independent of Bob's choice of measurement setting.

This can be taken to reflect the inability, under relativistic constraints, for information about Bob's settings to reach Alice's site in time to influence her outcomes.

Reasons to be doubtful:

- Probabilities refer to **ensembles**. The entire ensembles must be space-like separated from each other; and they may be unbounded in size . . .
- No-Signalling can be proved for ordinary QM with classical background. Indeed, it is a formal property of families of commuting observables, without any reference to causal background.

No-Signalling

The usual story:

When Alice and Bob perform measurements in spacelike-separated locations, the marginal probabilities for Alice's observations of outcomes for her measurements are independent of Bob's choice of measurement setting.

This can be taken to reflect the inability, under relativistic constraints, for information about Bob's settings to reach Alice's site in time to influence her outcomes.

Reasons to be doubtful:

- Probabilities refer to **ensembles**. The entire ensembles must be space-like separated from each other; and they may be unbounded in size . . .
- No-Signalling can be proved for ordinary QM with classical background. Indeed, it is a formal property of families of commuting observables, without any reference to causal background.

See e.g. the form of No-Signalling shown by S. Abramsky and A. Brandenburger, The sheaf-theoretic structure of non-locality and contextuality. *New Journal of Physics*, 13(2011):113036, 2011.

No-Signalling

The usual story:

When Alice and Bob perform measurements in spacelike-separated locations, the marginal probabilities for Alice's observations of outcomes for her measurements are independent of Bob's choice of measurement setting.

This can be taken to reflect the inability, under relativistic constraints, for information about Bob's settings to reach Alice's site in time to influence her outcomes.

Reasons to be doubtful:

- Probabilities refer to **ensembles**. The entire ensembles must be space-like separated from each other; and they may be unbounded in size . . .
- No-Signalling can be proved for ordinary QM with classical background. Indeed, it is a formal property of families of commuting observables, without any reference to causal background.

See e.g. the form of No-Signalling shown by S. Abramsky and A. Brandenburger, The sheaf-theoretic structure of non-locality and contextuality. *New Journal of Physics*, 13(2011):113036, 2011.

So the underlying structural reasons **why** it holds seem to be different.

Free Choice of Measurements

Free Choice of Measurements

This is a condition on hidden-variable theories, that the probability distribution on the hidden variable is independent of the probabilities on measurement settings induced by a given value of the hidden variable.

Free Choice of Measurements

This is a condition on hidden-variable theories, that the probability distribution on the hidden variable is independent of the probabilities on measurement settings induced by a given value of the hidden variable.

Otherwise, if for example the hidden variable actually **determines** which measurement settings can be chosen, then the hidden-variable theory becomes a self-fulfilling prophecy, and **any** empirical behaviour can trivially be reproduced.

Free Choice of Measurements

This is a condition on hidden-variable theories, that the probability distribution on the hidden variable is independent of the probabilities on measurement settings induced by a given value of the hidden variable.

Otherwise, if for example the hidden variable actually **determines** which measurement settings can be chosen, then the hidden-variable theory becomes a self-fulfilling prophecy, and **any** empirical behaviour can trivially be reproduced.

Without this condition, or something like it, it seems impossible to prove any substantive no-go theorem.

Free Choice of Measurements

This is a condition on hidden-variable theories, that the probability distribution on the hidden variable is independent of the probabilities on measurement settings induced by a given value of the hidden variable.

Otherwise, if for example the hidden variable actually **determines** which measurement settings can be chosen, then the hidden-variable theory becomes a self-fulfilling prophecy, and **any** empirical behaviour can trivially be reproduced.

Without this condition, or something like it, it seems impossible to prove any substantive no-go theorem.

Alternative name: λ -independence.

Free Choice of Measurements

This is a condition on hidden-variable theories, that the probability distribution on the hidden variable is independent of the probabilities on measurement settings induced by a given value of the hidden variable.

Otherwise, if for example the hidden variable actually **determines** which measurement settings can be chosen, then the hidden-variable theory becomes a self-fulfilling prophecy, and **any** empirical behaviour can trivially be reproduced.

Without this condition, or something like it, it seems impossible to prove any substantive no-go theorem.

Alternative name: λ -independence.

Apparently of a very different character to No-Signalling!

The Setting

The Setting

The Abramsky-Brandenburger sheaf-theoretic setting, providing a common generalisation of non-locality theory (Bell scenarios) and contextuality theory (Kochen-Specker configurations etc.)

The Setting

The Abramsky-Brandenburger sheaf-theoretic setting, providing a common generalisation of non-locality theory (Bell scenarios) and contextuality theory (Kochen-Specker configurations etc.)

Measurement cover: a pair (X, \mathcal{U}) , where X is a finite set, and \mathcal{U} is a family of subsets of X .

The Setting

The Abramsky-Brandenburger sheaf-theoretic setting, providing a common generalisation of non-locality theory (Bell scenarios) and contextuality theory (Kochen-Specker configurations etc.)

Measurement cover: a pair (X, \mathcal{U}) , where X is a finite set, and \mathcal{U} is a family of subsets of X .

Fix a set of **outcomes** O . $\Omega := O^X$ is a set of **canonical hidden variables**.

The Setting

The Abramsky-Brandenburger sheaf-theoretic setting, providing a common generalisation of non-locality theory (Bell scenarios) and contextuality theory (Kochen-Specker configurations etc.)

Measurement cover: a pair (X, \mathcal{U}) , where X is a finite set, and \mathcal{U} is a family of subsets of X .

Fix a set of **outcomes** O . $\Omega := O^X$ is a set of **canonical hidden variables**.

$D^\pm(S)$ is the set of **signed probability measures** on a (finite) set S , *i.e.* maps $d : S \rightarrow \mathbb{R}$ such that

$$\sum_{x \in S} d(x) = 1.$$

The Setting

The Abramsky-Brandenburger sheaf-theoretic setting, providing a common generalisation of non-locality theory (Bell scenarios) and contextuality theory (Kochen-Specker configurations etc.)

Measurement cover: a pair (X, \mathcal{U}) , where X is a finite set, and \mathcal{U} is a family of subsets of X .

Fix a set of **outcomes** O . $\Omega := O^X$ is a set of **canonical hidden variables**.

$D^\pm(S)$ is the set of **signed probability measures** on a (finite) set S , *i.e.* maps $d : S \rightarrow \mathbb{R}$ such that

$$\sum_{x \in S} d(x) = 1.$$

$D^+(S)$ is the subset of $D^\pm(S)$ of *bona fide* probability measures, *i.e.* valued in the non-negative reals.

The Setting

The Abramsky-Brandenburger sheaf-theoretic setting, providing a common generalisation of non-locality theory (Bell scenarios) and contextuality theory (Kochen-Specker configurations etc.)

Measurement cover: a pair (X, \mathcal{U}) , where X is a finite set, and \mathcal{U} is a family of subsets of X .

Fix a set of **outcomes** O . $\Omega := O^X$ is a set of **canonical hidden variables**.

$D^\pm(S)$ is the set of **signed probability measures** on a (finite) set S , *i.e.* maps $d : S \rightarrow \mathbb{R}$ such that

$$\sum_{x \in S} d(x) = 1.$$

$D^+(S)$ is the subset of $D^\pm(S)$ of *bona fide* probability measures, *i.e.* valued in the non-negative reals.

Given (X, \mathcal{U}) and O , we define a set of **atomic events** E :

$$E := \{(U, s) \mid U \in \mathcal{U} \wedge s \in O^U\}.$$

The Setting

The Abramsky-Brandenburger sheaf-theoretic setting, providing a common generalisation of non-locality theory (Bell scenarios) and contextuality theory (Kochen-Specker configurations etc.)

Measurement cover: a pair (X, \mathcal{U}) , where X is a finite set, and \mathcal{U} is a family of subsets of X .

Fix a set of **outcomes** O . $\Omega := O^X$ is a set of **canonical hidden variables**.

$D^\pm(S)$ is the set of **signed probability measures** on a (finite) set S , *i.e.* maps $d : S \rightarrow \mathbb{R}$ such that

$$\sum_{x \in S} d(x) = 1.$$

$D^+(S)$ is the subset of $D^\pm(S)$ of *bona fide* probability measures, *i.e.* valued in the non-negative reals.

Given (X, \mathcal{U}) and O , we define a set of **atomic events** E :

$$E := \{(U, s) \mid U \in \mathcal{U} \wedge s \in O^U\}.$$

Thus (U, s) is the event that the measurements in U were performed, and the outcome $s(x)$ was observed for each $x \in U$.

Empirical Models

Empirical Models

An **empirical model** over $((X, \mathcal{U}), O)$ is a probability distribution $e \in D^+(E)$ such that, for each $U \in \mathcal{U}$,

$$e(U) := \sum_{s \in O^U} e(U, s) > 0.$$

Empirical Models

An **empirical model** over $((X, \mathcal{U}), O)$ is a probability distribution $e \in D^+(E)$ such that, for each $U \in \mathcal{U}$,

$$e(U) := \sum_{s \in O^U} e(U, s) > 0.$$

For each $U \in \mathcal{U}$, e determines a probability distribution $e_U \in D^+(O^U)$ as the conditional probability $e_U(s) := e((U, s)|U)$.

Empirical Models

An **empirical model** over $((X, \mathcal{U}), O)$ is a probability distribution $e \in D^+(E)$ such that, for each $U \in \mathcal{U}$,

$$e(U) := \sum_{s \in O^U} e(U, s) > 0.$$

For each $U \in \mathcal{U}$, e determines a probability distribution $e_U \in D^+(O^U)$ as the conditional probability $e_U(s) := e((U, s)|U)$.

Empirical models generalise the standard probability tables familiar from non-locality theory.

Empirical Models

An **empirical model** over $((X, \mathcal{U}), O)$ is a probability distribution $e \in D^+(E)$ such that, for each $U \in \mathcal{U}$,

$$e(U) := \sum_{s \in O^U} e(U, s) > 0.$$

For each $U \in \mathcal{U}$, e determines a probability distribution $e_U \in D^+(O^U)$ as the conditional probability $e_U(s) := e((U, s)|U)$.

Empirical models generalise the standard probability tables familiar from non-locality theory.

E.g. take $X = A \sqcup B$ to be Alice's measurements plus Bob's measurements, and

$$\mathcal{U} := \{\{a, b\} \mid a \in A, b \in B\}.$$

Empirical Models

An **empirical model** over $((X, \mathcal{U}), O)$ is a probability distribution $e \in D^+(E)$ such that, for each $U \in \mathcal{U}$,

$$e(U) := \sum_{s \in O^U} e(U, s) > 0.$$

For each $U \in \mathcal{U}$, e determines a probability distribution $e_U \in D^+(O^U)$ as the conditional probability $e_U(s) := e((U, s)|U)$.

Empirical models generalise the standard probability tables familiar from non-locality theory.

E.g. take $X = A \sqcup B$ to be Alice's measurements plus Bob's measurements, and

$$\mathcal{U} := \{\{a, b\} \mid a \in A, b \in B\}.$$

Also generalises K-S configurations: given a family of observables X , take \mathcal{U} to be the family of commuting subsets.

Hidden-Variable Models

Hidden-Variable Models

A **signed canonical hidden-variable model** (schv model) is a signed measure $m \in D^\pm(\Omega \times \mathcal{U})$.

Hidden-Variable Models

A **signed canonical hidden-variable model** (schv model) is a signed measure $m \in D^\pm(\Omega \times \mathcal{U})$.

A **probabilistic canonical hidden-variable model** (pchv model) is an schv model p such that $p \in D^+(\Omega \times \mathcal{U})$.

Hidden-Variable Models

A **signed canonical hidden-variable model** (schv model) is a signed measure $m \in D^\pm(\Omega \times \mathcal{U})$.

A **probabilistic canonical hidden-variable model** (pchv model) is an schv model p such that $p \in D^+(\Omega \times \mathcal{U})$.

Given $(U, s) \in E$, we define

$$\Omega(U, s) := \{\omega \in \Omega \mid \omega|_U = s\}.$$

Hidden-Variable Models

A **signed canonical hidden-variable model** (schv model) is a signed measure $m \in D^\pm(\Omega \times \mathcal{U})$.

A **probabilistic canonical hidden-variable model** (pchv model) is an schv model p such that $p \in D^+(\Omega \times \mathcal{U})$.

Given $(U, s) \in E$, we define

$$\Omega(U, s) := \{\omega \in \Omega \mid \omega|_U = s\}.$$

This is the set of canonical hidden variables which are consistent with the atomic event (U, s) .

Hidden-Variable Models

A **signed canonical hidden-variable model** (schv model) is a signed measure $m \in D^\pm(\Omega \times \mathcal{U})$.

A **probabilistic canonical hidden-variable model** (pchv model) is an schv model p such that $p \in D^+(\Omega \times \mathcal{U})$.

Given $(U, s) \in E$, we define

$$\Omega(U, s) := \{\omega \in \Omega \mid \omega|_U = s\}.$$

This is the set of canonical hidden variables which are consistent with the atomic event (U, s) .

Note that, for each $U \in \mathcal{U}$, the sets $\Omega(U, s)$ as s ranges over O^U partition Ω .

Hidden-Variable Models

A **signed canonical hidden-variable model** (schv model) is a signed measure $m \in D^\pm(\Omega \times \mathcal{U})$.

A **probabilistic canonical hidden-variable model** (pchv model) is an schv model p such that $p \in D^+(\Omega \times \mathcal{U})$.

Given $(U, s) \in E$, we define

$$\Omega(U, s) := \{\omega \in \Omega \mid \omega|_U = s\}.$$

This is the set of canonical hidden variables which are consistent with the atomic event (U, s) .

Note that, for each $U \in \mathcal{U}$, the sets $\Omega(U, s)$ as s ranges over O^U partition Ω .

An schv model m determines a signed measure $\hat{m} \in D^\pm(E)$ by marginalization:

$$\hat{m}(U, s) := \sum_{\omega \in \Omega(U, s)} m(\omega, U).$$

Hidden-Variable Models

A **signed canonical hidden-variable model** (schv model) is a signed measure $m \in D^\pm(\Omega \times \mathcal{U})$.

A **probabilistic canonical hidden-variable model** (pchv model) is an schv model p such that $p \in D^+(\Omega \times \mathcal{U})$.

Given $(U, s) \in E$, we define

$$\Omega(U, s) := \{\omega \in \Omega \mid \omega|_U = s\}.$$

This is the set of canonical hidden variables which are consistent with the atomic event (U, s) .

Note that, for each $U \in \mathcal{U}$, the sets $\Omega(U, s)$ as s ranges over O^U partition Ω .

An schv model m determines a signed measure $\hat{m} \in D^\pm(E)$ by marginalization:

$$\hat{m}(U, s) := \sum_{\omega \in \Omega(U, s)} m(\omega, U).$$

We say that an schv model m **realizes** an empirical model e if $\hat{m} = e$.

The need for the FCM assumption

The need for the FCM assumption

Proposition

For every empirical model e , there is a pchv model p which realizes e .

The need for the FCM assumption

Proposition

For every empirical model e , there is a pchv model p which realizes e .

So without additional constraints realization by deterministic hidden variables is trivially achieved.

The need for the FCM assumption

Proposition

For every empirical model e , there is a pchv model p which realizes e .

So without additional constraints realization by deterministic hidden variables is trivially achieved.

The key condition is **free choice of measurements** (FCM, aka Lambda Independence):

The distribution on the hidden variables should be statistically independent of the choice of measurement context.

The need for the FCM assumption

Proposition

For every empirical model e , there is a pchv model p which realizes e .

So without additional constraints realization by deterministic hidden variables is trivially achieved.

The key condition is **free choice of measurements** (FCM, aka Lambda Independence):

The distribution on the hidden variables should be statistically independent of the choice of measurement context.

Formally, an schv model m satisfies FCM if it factors as a product: $m = m_{\Omega}m_{\mathcal{U}}$, where $m_{\Omega} \in D^{\pm}(\Omega)$ and $m_{\mathcal{U}} \in D^{\pm}(\mathcal{U})$ are the marginals of m .

The need for the FCM assumption

Proposition

For every empirical model e , there is a schv model p which realizes e .

So without additional constraints realization by deterministic hidden variables is trivially achieved.

The key condition is **free choice of measurements** (FCM, aka Lambda Independence):

The distribution on the hidden variables should be statistically independent of the choice of measurement context.

Formally, an schv model m satisfies FCM if it factors as a product: $m = m_{\Omega} m_{\mathcal{U}}$, where $m_{\Omega} \in D^{\pm}(\Omega)$ and $m_{\mathcal{U}} \in D^{\pm}(\mathcal{U})$ are the marginals of m .

Proposition

An schv model m satisfies FCM iff for all $\omega \in \Omega$, $U, U' \in \mathcal{U}$,

$$m(\omega|U) = m(\omega|U').$$

Bell in Context

Bell in Context

We say that an empirical model which is realized by a pchv model satisfying FCM admits **local hidden variables**.

Bell in Context

We say that an empirical model which is realized by a pchv model satisfying FCM admits **local hidden variables**.

This definition is equivalent to the standard definitions, which allow a broader class of hidden variable models.

Bell in Context

We say that an empirical model which is realized by a pchv model satisfying FCM admits **local hidden variables**.

This definition is equivalent to the standard definitions, which allow a broader class of hidden variable models.

Theorem (Bell's Theorem)

There are empirical models which can be realized in quantum mechanics which do not admit local hidden variables.

No-Signalling

No-Signalling

An empirical model e satisfies **No-Signalling** if for all $U, U' \in \mathcal{U}$, $s \in \mathcal{O}^{U \cap U'}$:

$$e_U(s) = e_{U'}(s).$$

No-Signalling

An empirical model e satisfies **No-Signalling** if for all $U, U' \in \mathcal{U}$, $s \in \mathcal{O}^{U \cap U'}$:

$$e_U(s) = e_{U'}(s).$$

This says that the distributions conditioned on different measurement choices have common marginals.

No-Signalling

An empirical model e satisfies **No-Signalling** if for all $U, U' \in \mathcal{U}$, $s \in \mathcal{O}^{U \cap U'}$:

$$e_U(s) = e_{U'}(s).$$

This says that the distributions conditioned on different measurement choices have common marginals.

Thus the choice of additional measurements $U \setminus V$ outside a compatible set V has no effect on the observed statistics for the measurement outcomes in V .

No-Signalling

An empirical model e satisfies **No-Signalling** if for all $U, U' \in \mathcal{U}$, $s \in \mathcal{O}^{U \cap U'}$:

$$e_U(s) = e_{U'}(s).$$

This says that the distributions conditioned on different measurement choices have common marginals.

Thus the choice of additional measurements $U \setminus V$ outside a compatible set V has no effect on the observed statistics for the measurement outcomes in V .

This is easily seen to be equivalent to the standard formulation of No-Signalling in Bell-type scenarios, and to be satisfied generally in quantum mechanics (Abramsky-Brandenburger 2011).

No-Signalling

An empirical model e satisfies **No-Signalling** if for all $U, U' \in \mathcal{U}$, $s \in \mathcal{O}^{U \cap U'}$:

$$e_U(s) = e_{U'}(s).$$

This says that the distributions conditioned on different measurement choices have common marginals.

Thus the choice of additional measurements $U \setminus V$ outside a compatible set V has no effect on the observed statistics for the measurement outcomes in V .

This is easily seen to be equivalent to the standard formulation of No-Signalling in Bell-type scenarios, and to be satisfied generally in quantum mechanics (Abramsky-Brandenburger 2011).

Proposition

Let m be an schv model satisfying FCM which realizes an empirical model $e = \hat{m}$. Then e satisfies No-Signalling.

No-Signalling

An empirical model e satisfies **No-Signalling** if for all $U, U' \in \mathcal{U}$, $s \in \mathcal{O}^{U \cap U'}$:

$$e_U(s) = e_{U'}(s).$$

This says that the distributions conditioned on different measurement choices have common marginals.

Thus the choice of additional measurements $U \setminus V$ outside a compatible set V has no effect on the observed statistics for the measurement outcomes in V .

This is easily seen to be equivalent to the standard formulation of No-Signalling in Bell-type scenarios, and to be satisfied generally in quantum mechanics (Abramsky-Brandenburger 2011).

Proposition

Let m be an schv model satisfying FCM which realizes an empirical model $e = \hat{m}$. Then e satisfies No-Signalling.

Note however that the converse does not hold, even for pchv models.

The Main Result

We now ask:

*Which empirical models can be realised by local hidden variables, **if we allow signed measures (negative probabilities) on the hidden variables?***

The Main Result

We now ask:

*Which empirical models can be realised by local hidden variables, **if we allow signed measures (negative probabilities) on the hidden variables?***

It is important to note that the empirical models are *bona fide* probabilistic models. The negative probabilities are used only on the hidden variables, and must cancel out to yield non-negative probabilities on the observed quantities.

The Main Result

We now ask:

*Which empirical models can be realised by local hidden variables, if we **allow signed measures (negative probabilities)** on the hidden variables?*

It is important to note that the empirical models are *bona fide* probabilistic models. The negative probabilities are used only on the hidden variables, and must cancel out to yield non-negative probabilities on the observed quantities.

Negative probabilities have been considered by Wigner, Dirac, Feynman, . . .

The Main Result

We now ask:

*Which empirical models can be realised by local hidden variables, **if we allow signed measures (negative probabilities) on the hidden variables?***

It is important to note that the empirical models are *bona fide* probabilistic models. The negative probabilities are used only on the hidden variables, and must cancel out to yield non-negative probabilities on the observed quantities.

Negative probabilities have been considered by Wigner, Dirac, Feynman, . . .

Feynman:

The only difference between a probabilistic classical world and the equations of the quantum world is that somehow or other it appears as if the probabilities would have to go negative . . .

The Main Result

We now ask:

*Which empirical models can be realised by local hidden variables, if we **allow signed measures (negative probabilities)** on the hidden variables?*

It is important to note that the empirical models are *bona fide* probabilistic models. The negative probabilities are used only on the hidden variables, and must cancel out to yield non-negative probabilities on the observed quantities.

Negative probabilities have been considered by Wigner, Dirac, Feynman, . . .

Feynman:

The only difference between a probabilistic classical world and the equations of the quantum world is that somehow or other it appears as if the probabilities would have to go negative . . .

Theorem

Empirical models have local hidden-variable realizations with negative probabilities if and only if they satisfy no-signalling.

Linear Span Theorem

Linear Span Theorem

The fact that all no-signalling empirical models admit local hidden variables with signed measures is a consequence of the following:

Linear Span Theorem

The fact that all no-signalling empirical models admit local hidden variables with signed measures is a consequence of the following:

Theorem

The linear subspace generated by the local models over an arbitrary measurement cover \mathcal{U} coincides with that generated by the no-signalling models. Their common dimension is

$$D := \sum_{U \in \Sigma} (I - 1)^{|U|}$$

where $I = |O|$ and Σ is the abstract simplicial complex generated by \mathcal{U} .

Linear Span Theorem

The fact that all no-signalling empirical models admit local hidden variables with signed measures is a consequence of the following:

Theorem

The linear subspace generated by the local models over an arbitrary measurement cover \mathcal{U} coincides with that generated by the no-signalling models. Their common dimension is

$$D := \sum_{U \in \Sigma} (I - 1)^{|U|}$$

where $I = |O|$ and Σ is the abstract simplicial complex generated by \mathcal{U} .

Since the local models are included in the no-signalling models, this is proved by showing that every compatible model is determined by linear equations in D variables; while there are D linearly independent local models.

Example: PR Boxes have global sections over \mathbb{R}

Example: PR Boxes have global sections over \mathbb{R}

The 'Popescu-Rohrlich box':

	(0, 0)	(1, 0)	(0, 1)	(1, 1)
(a, b)	1/2	0	0	1/2
(a', b)	1/2	0	0	1/2
(a, b')	1/2	0	0	1/2
(a', b')	0	1/2	1/2	0

Example: PR Boxes have global sections over \mathbb{R}

The 'Popescu-Rohrlich box':

	(0, 0)	(1, 0)	(0, 1)	(1, 1)
(a, b)	1/2	0	0	1/2
(a', b)	1/2	0	0	1/2
(a, b')	1/2	0	0	1/2
(a', b')	0	1/2	1/2	0

The PR boxes exhibit super-quantum correlations, and cannot be realized in quantum mechanics.

Example: PR Boxes have global sections over \mathbb{R}

The 'Popescu-Rohrlich box':

	(0, 0)	(1, 0)	(0, 1)	(1, 1)
(a, b)	1/2	0	0	1/2
(a', b)	1/2	0	0	1/2
(a, b')	1/2	0	0	1/2
(a', b')	0	1/2	1/2	0

The PR boxes exhibit super-quantum correlations, and cannot be realized in quantum mechanics.

Example solution:

$$[1/2, 0, 0, 0, -1/2, 0, 1/2, 0, -1/2, 1/2, 0, 0, 1/2, 0, 0, 0].$$

Example: PR Boxes have global sections over \mathbb{R}

The 'Popescu-Rohrlich box':

	(0, 0)	(1, 0)	(0, 1)	(1, 1)
(a, b)	1/2	0	0	1/2
(a', b)	1/2	0	0	1/2
(a, b')	1/2	0	0	1/2
(a', b')	0	1/2	1/2	0

The PR boxes exhibit super-quantum correlations, and cannot be realized in quantum mechanics.

Example solution:

$$[1/2, 0, 0, 0, -1/2, 0, 1/2, 0, -1/2, 1/2, 0, 0, 1/2, 0, 0, 0].$$

This vector can be taken as giving a **local hidden-variable realization of the PR box using negative probabilities**. Similar explicit realizations can be given for the other PR boxes.

The Equivalence

The Equivalence

The equivalence between No-Signalling and FCM is stated precisely in the following results:

The Equivalence

The equivalence between No-Signalling and FCM is stated precisely in the following results:

Proposition

Let m be an schv model satisfying FCM which realizes an empirical model $e = \hat{m}$. Then e satisfies No-Signalling.

The Equivalence

The equivalence between No-Signalling and FCM is stated precisely in the following results:

Proposition

Let m be an schv model satisfying FCM which realizes an empirical model $e = \hat{m}$. Then e satisfies No-Signalling.

As a consequence of our main theorem:

Proposition

Every empirical model e satisfying No-Signalling is realized by some schv model satisfying FCM.

The Equivalence

The equivalence between No-Signalling and FCM is stated precisely in the following results:

Proposition

Let m be an schv model satisfying FCM which realizes an empirical model $e = \hat{m}$. Then e satisfies No-Signalling.

As a consequence of our main theorem:

Proposition

Every empirical model e satisfying No-Signalling is realized by some schv model satisfying FCM.

Hence we obtain the equivalence:

Theorem

An empirical model is No-Signalling if and only if it is realized by an schv model satisfying FCM.

But what are these negative probabilities?

But what are these negative probabilities?

We shall sketch an 'operational interpretation' of negative probabilities, and how this yields a way of realising arbitrary no-signalling devices.

But what are these negative probabilities?

We shall sketch an 'operational interpretation' of negative probabilities, and how this yields a way of realising arbitrary no-signalling devices.

Not 'physical'. Perhaps best thought of as a simulation.

But what are these negative probabilities?

We shall sketch an 'operational interpretation' of negative probabilities, and how this yields a way of realising arbitrary no-signalling devices.

Not 'physical'. Perhaps best thought of as a simulation.

The basic idea: 'push the minus signs inwards'. We take **signed events**.

But what are these negative probabilities?

We shall sketch an 'operational interpretation' of negative probabilities, and how this yields a way of realising arbitrary no-signalling devices.

Not 'physical'. Perhaps best thought of as a simulation.

The basic idea: 'push the minus signs inwards'. We take **signed events**.

The idea (in a frequentist setting) is that events of opposite sign cancel.

$$\bullet^+, \bullet^- \longrightarrow$$

But what are these negative probabilities?

We shall sketch an 'operational interpretation' of negative probabilities, and how this yields a way of realising arbitrary no-signalling devices.

Not 'physical'. Perhaps best thought of as a simulation.

The basic idea: 'push the minus signs inwards'. We take **signed events**.

The idea (in a frequentist setting) is that events of opposite sign cancel.

$$\bullet^+, \bullet^- \longrightarrow$$

We also have 'creation' as well as 'annihilation', since adding opposite events will not affect the **signed relative frequencies**:

$$\longrightarrow \bullet^+, \bullet^-$$

But what are these negative probabilities?

We shall sketch an 'operational interpretation' of negative probabilities, and how this yields a way of realising arbitrary no-signalling devices.

Not 'physical'. Perhaps best thought of as a simulation.

The basic idea: 'push the minus signs inwards'. We take **signed events**.

The idea (in a frequentist setting) is that events of opposite sign cancel.

$$\bullet^+, \bullet^- \longrightarrow$$

We also have 'creation' as well as 'annihilation', since adding opposite events will not affect the **signed relative frequencies**:

$$\longrightarrow \bullet^+, \bullet^-$$

The signed relative frequency of an event \bullet in an ensemble of size N will be

$$\frac{n^+ - n^-}{N}$$

where \bullet^+ occurs n^+ times and \bullet^- occurs n^- times in the ensemble.

More formally

More formally

There is a map

$$\theta_x : D^\pm(X) \longrightarrow D^+(X + X)$$

which puts positive weight $d(x) \geq 0$ on x on positive copy (first summand) and (absolute value of) negative weight on second copy.

More formally

There is a map

$$\theta_X : D^\pm(X) \longrightarrow D^+(X + X)$$

which puts positive weight $d(x) \geq 0$ on x on positive copy (first summand) and (absolute value of) negative weight on second copy.

There is a map

$$\eta_X : D^+(X + X) \longrightarrow D^\pm(X)$$

which sets $\eta_X(d)(x) := d(x^+) - d(x^-)$.

More formally

There is a map

$$\theta_X : D^\pm(X) \longrightarrow D^+(X + X)$$

which puts positive weight $d(x) \geq 0$ on x on positive copy (first summand) and (absolute value of) negative weight on second copy.

There is a map

$$\eta_X : D^+(X + X) \longrightarrow D^\pm(X)$$

which sets $\eta_X(d)(x) := d(x^+) - d(x^-)$.

Caveats:

More formally

There is a map

$$\theta_X : D^\pm(X) \longrightarrow D^+(X + X)$$

which puts positive weight $d(x) \geq 0$ on x on positive copy (first summand) and (absolute value of) negative weight on second copy.

There is a map

$$\eta_X : D^+(X + X) \longrightarrow D^\pm(X)$$

which sets $\eta_X(d)(x) := d(x^+) - d(x^-)$.

Caveats:

- Need to renormalise $\theta_X(d)$.

More formally

There is a map

$$\theta_X : D^\pm(X) \longrightarrow D^+(X + X)$$

which puts positive weight $d(x) \geq 0$ on x on positive copy (first summand) and (absolute value of) negative weight on second copy.

There is a map

$$\eta_X : D^+(X + X) \longrightarrow D^\pm(X)$$

which sets $\eta_X(d)(x) := d(x^+) - d(x^-)$.

Caveats:

- Need to renormalise $\theta_X(d)$.
- Naturality: need to work with reflection to normalised subcategory in which positive and negative weights have been cancelled.

More formally

There is a map

$$\theta_X : D^\pm(X) \longrightarrow D^+(X + X)$$

which puts positive weight $d(x) \geq 0$ on x on positive copy (first summand) and (absolute value of) negative weight on second copy.

There is a map

$$\eta_X : D^+(X + X) \longrightarrow D^\pm(X)$$

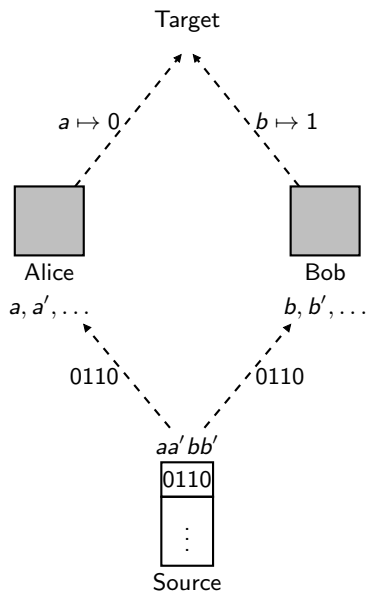
which sets $\eta_X(d)(x) := d(x^+) - d(x^-)$.

Caveats:

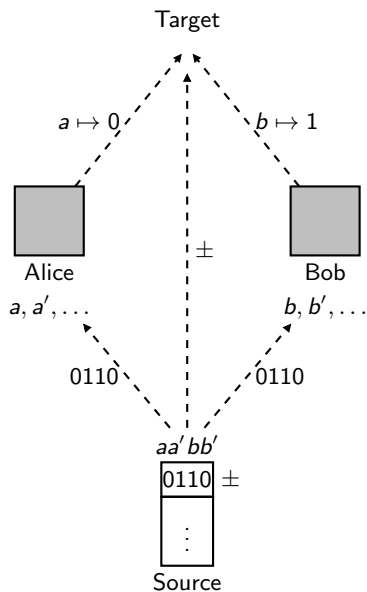
- Need to renormalise $\theta_X(d)$.
- Naturality: need to work with reflection to normalised subcategory in which positive and negative weights have been cancelled.

This is a (pale) reflection of the Hahn-Jordan decomposition.

The Mermin instruction set picture



Signed version of the Mermin instruction set picture



Final Remarks

Final Remarks

So we have an 'operational semantics' for simulating arbitrary no-signalling devices.

Final Remarks

So we have an 'operational semantics' for simulating arbitrary no-signalling devices.

Rather than taking e.g. PR boxes as 'black boxes' we have some story about how they might work inside.

Final Remarks

So we have an 'operational semantics' for simulating arbitrary no-signalling devices.

Rather than taking e.g. PR boxes as 'black boxes' we have some story about how they might work inside.

- Uses?

Final Remarks

So we have an 'operational semantics' for simulating arbitrary no-signalling devices.

Rather than taking e.g. PR boxes as 'black boxes' we have some story about how they might work inside.

- Uses?
- Does 'cancellation' imply or otherwise relate to **retrocausality**?